

Libra/CP2K Interface

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CP2K

- CP2K is a quantum chemistry and solid state physics software package that can perform atomistic simulations of solid state, liquid, molecular, periodic, material, crystal, and biological systems.
- CP2K provides a general framework for different modeling methods such as DFT using the mixed Gaussian and plane waves approaches GPW and GAPW.
- Supported theory levels include DFTB, LDA, GGA, MP2, RPA, semi-empirical methods (AM1, PM3, PM6, RM1, MNDO, ...), and classical force fields (AMBER, CHARMM, ...).
- CP2K can do simulations of molecular dynamics, metadynamics, Monte Carlo, Ehrenfest dynamics, vibrational analysis, core level spectroscopy, energy minimization, and transition state optimization using NEB or dimer method.

Theory

- A molecular orbital (MO) is defined as a linear combination of atomic orbitals:

$$|\Psi_n\rangle = \sum_i c_i |\phi_i\rangle$$

- The MO overlap:

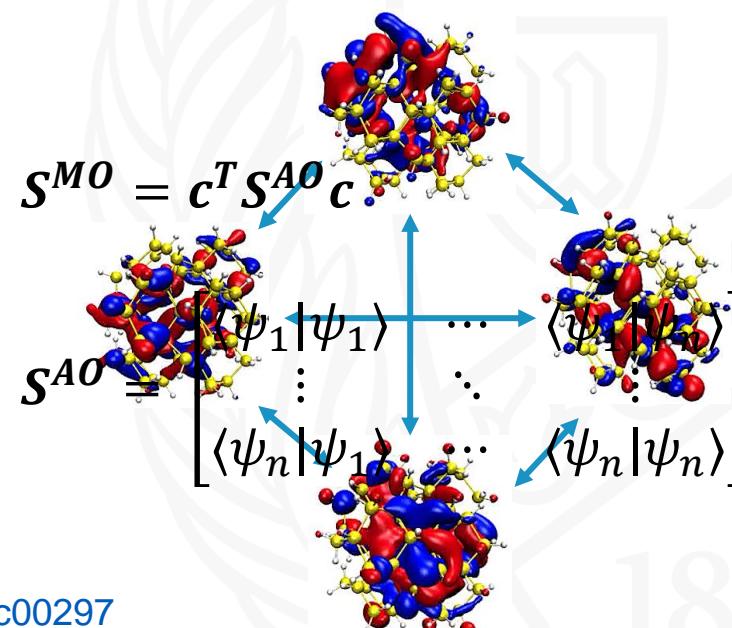
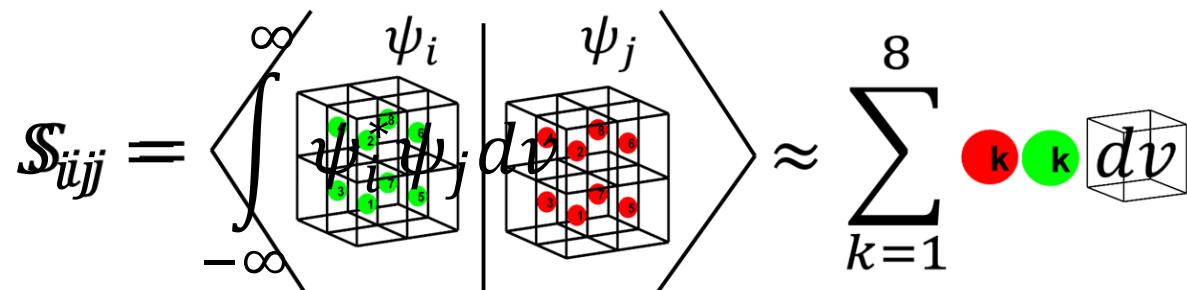
$$S_{nm} = \langle \Psi_n | \Psi_m \rangle = \sum_{i,j} c_i^* c_j \langle \phi_i | \phi_j \rangle$$

the $\langle \phi_i | \phi_j \rangle$ is the atomic orbital overlap. So the MO overlap matrix (S^{MO}) can be written as follows where c is the matrix of molecular orbital coefficients and S^{AO} is the atomic orbital overlap matrix:

$$S^{MO} = c^T S^{AO} c$$

Molecular orbital integrals

- Grid-based approach using *.cube* files
 - Easy to implement
 - Most codes can output these file
 - Not suitable for large structures with large number of states
- Analytical approach
 - Suitable for large systems and large number of states
 - One interface cannot directly be used for many codes



Smith, Shakiba, Akimov, J. Chem. Theory Comp. 2021, 17, 678–693

Shakiba, Stippel, Akimov, J. Chem. Theory Comp. 2022, DOI: [10.1021/acs.jctc.2c00297](https://doi.org/10.1021/acs.jctc.2c00297)

Gaussian type orbitals (GTO)

- Atom centered basis sets

$$\varphi(r - R; n, \xi) = N(x - R_x)^{n_x} (y - R_y)^{n_y} (z - R_z)^{n_z} \times \exp(-\xi(r - R)^2)$$

$$N = \left(\frac{2\xi}{\pi}\right)^{\frac{3}{4}} (4\xi)^{(n_x+n_y+n_z)/2} \times \left((2n_x - 1)!! (2n_y - 1)!! (2n_z - 1)!!\right)^{-\frac{1}{2}}$$

$$\Phi_i(\vec{r}) = \varphi(r; n_i, \xi) \cdot Y_{n_i, m_i}(\theta, \phi)$$

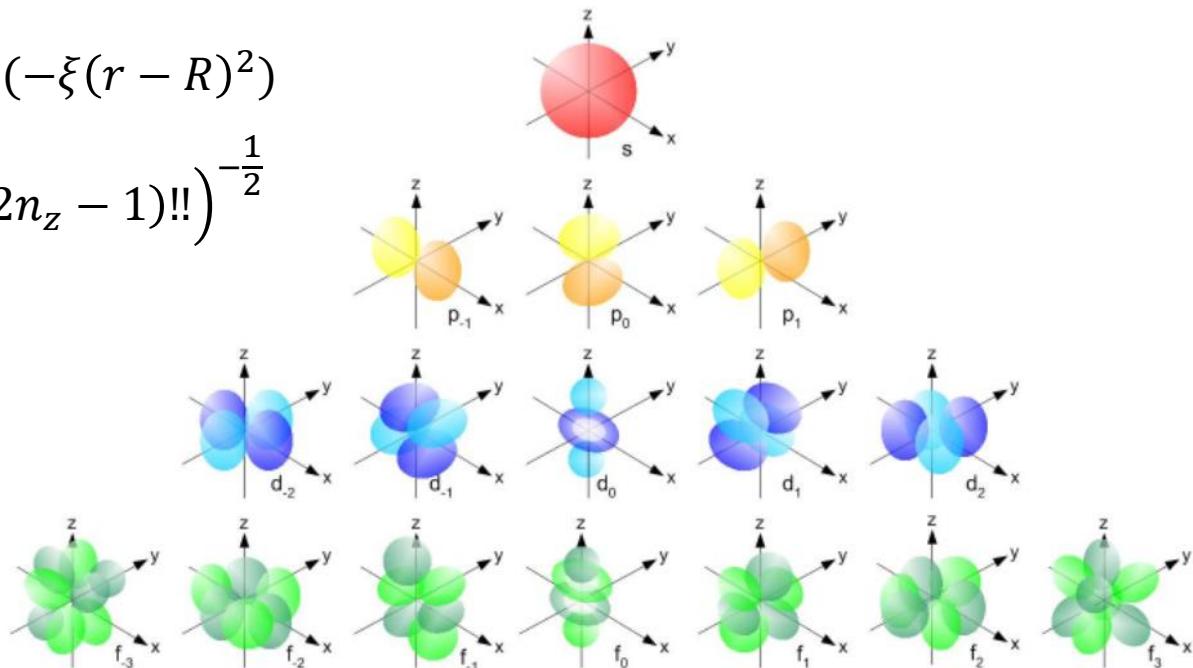
n shows the angular momentum value:

s-orbital $n = 0$

p-orbital $n = 1$

d-orbital $n = 2$

f-orbital $n = 3$



Recurrence relations between GTOs

$$\langle s|s \rangle = \left(\frac{\pi}{\xi}\right)^{\frac{3}{2}} \exp(-\xi(A - B)^2)$$

$$\langle p_i | s \rangle = (P_i - A_i) \langle s | s \rangle$$

$$\langle p_i | p_j \rangle = (P_j - B_j) \langle p_i | s \rangle + \frac{\delta_{ij}}{2\xi} \langle s | s \rangle$$

$$\langle d_{ij} | s \rangle = (P_j - A_j) \langle p_i | s \rangle + \frac{\delta_{ij}}{2\xi} \langle s | s \rangle$$

$$\langle d_{ij} | p_k \rangle = (P_k - B_k) \langle d_{ij} | s \rangle + \frac{\delta_{ik}}{2\xi} \langle p_j | s \rangle + \frac{\delta_{jk}}{2\xi} \langle p_i | s \rangle$$

$$\langle d_{ij} | d_{kl} \rangle = (P_l - B_l) \langle d_{ij} | p_k \rangle + \frac{\delta_{il}}{2\xi} \langle p_j | p_k \rangle + \frac{\delta_{jl}}{2\xi} \langle p_i | p_k \rangle + \frac{\delta_{kl}}{2\xi} \langle d_{ij} | s \rangle$$

$i, j, k, l = x, y, z$

Obara, and Saika, J. Chem. Phys. 1986, 84, 3963-3974.

Obara, and Saika, J. Chem. Phys. 1988, 89, 1540-1559.

Libint, Version 2.6.0 Edward F. Valeev, <http://libint.valeev.net>.



Open-source and OpenMP
parallelized code for computing
integrals between GTOs

1846

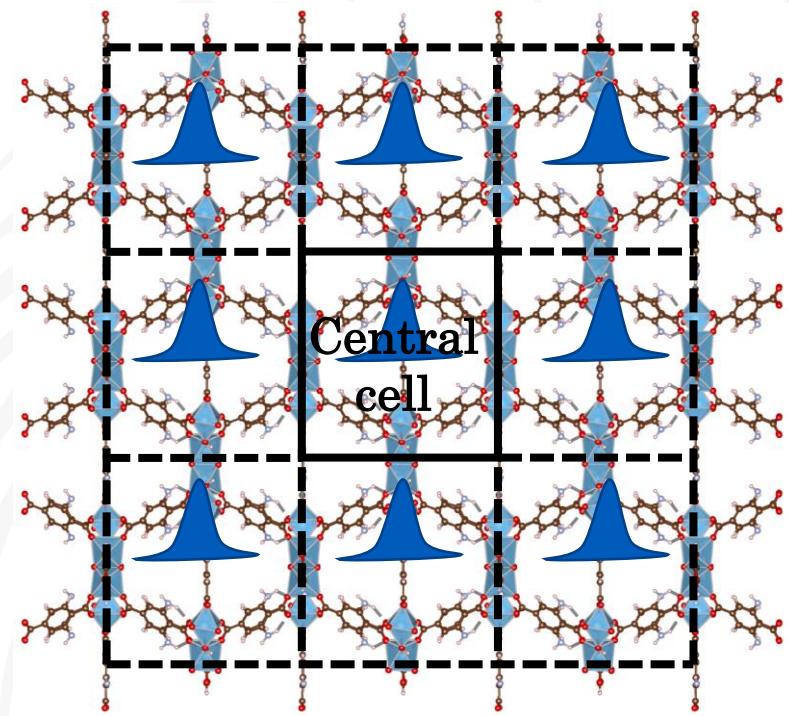
Periodic basis functions

- The Bloch function for K-point in a periodic structure is defined as:

$$\beta_a^k(r) = \frac{1}{\sqrt{N}} \sum_R \varphi_a(r - R) e^{ikR}$$

- Overlaps between Bloch functions of two different K-points:

$$S_{a,b}^k = \langle \beta_a^k | \beta_b^{k'} \rangle = \frac{1}{N} \int dr \sum_{R,R'} e^{-ikR} \varphi_a^*(r - R) e^{ik'R'} \varphi_b(r - R') \\ = \frac{1}{N} \int dr \sum_{R,R'} e^{i(kR - k'R')} \varphi_a^*(r - R) \varphi_b(r - R')$$



Aradi, Hourahine, Frauenheim, J. Phys. Chem. A 2007, 111, 5678–5684

Shakiba, Stippell, Li, Akimov, J. Chem. Theory Compu. 2022, DOI: [10.1021/acs.jctc.2c00297](https://doi.org/10.1021/acs.jctc.2c00297)

How Libra computes the MO overlap?

- Libra uses ‘*molden*’ file formats to read the molecular orbital coefficients, energies, occupation, spin, atomic coordinates, and basis set although it can use other file formats as well including *MOLog* printed out by CP2K.
- Here, we’ll show an example of the inner functions that compute the MO overlap matrix for one geometry. Note that the workflow is not like this and you only need to specify a couple of variables. The following shows how one can work with the Libra functions, for example, if one intends to write an on-the-fly computation of the NACs.

```

import os
import numpy as np
import matplotlib.pyplot as plt
from liblibra_core import *
from libra_py import CP2K_methods,molden_methods, data_conv, units
# Molden file name
molden_file_name = 'test.molden'
# Number of processors
nprocs = 16
# Spherical or Cartesian coordiantes? Spherical!
is_spherical = True
# The integration shells and angular momentum values
shell_1, l_vals =
molden_methods.molden_file_to_libint_shell(molden_file_name,
is_spherical)
# All of the eigenvectors and energies of the system
eig_vect_1, energies_1 =
molden_methods.eigenvalues_molden(molden_file_name, nbasis(shell_1),
l_vals)

```

```

# Resorting the molden indices
new_indices = CP2K_methods.resort_molog_eigenvalues(l_vals)
eigenvectors_1 = []
for j in range(len(eig_vect_1)):
    # the new and sorted eigenvector
    eigenvector_1 = eig_vect_1[j]
    eigenvector_1 = eigenvector_1[new_indices]
    # append it to the eigenvectors list
    eigenvectors_1.append(eigenvector_1)
eigenvectors_1 = np.array(eigenvectors_1)

```

```

# Alpha and Beta spin eigenvectors
# alpha -> even indices
alpha_eig_vects = eigenvectors_1[0::2]
alpha_eig_vals = energies_1[0::2]
# beta -> odd indices
beta_eig_vects = eigenvectors_1[1::2]
beta_eig_vals = energies_1[1::2]
# Compute AO overlap matrix
AO_S = compute_overlaps(shell_1, shell_1, nprocs)
# Converting to numpy array
AO_S = data_conv.MATRIX2numpy(AO_S)

print('The shape of the AO matrix... \n', AO_S.shape)
S_alpha = np.linalg.multi_dot([alpha_eig_vects, AO_S,
alpha_eig_vects.T])
S_beta = np.linalg.multi_dot([beta_eig_vects, AO_S,
beta_eig_vects.T])

```

Periodic calculations

```
AO_S = compute_overlaps(shell_1,shell_1,nprocs)
if is_periodic:
    cell = []
    cell.append(params['A_cell_vector'])
    cell.append(params['B_cell_vector'])
    cell.append(params['C_cell_vector'])
    cell = np.array(cell)*units.Angst
    # Generating translational vectors
    translational_vectors = params['translational_vectors']
    for i1 in range(len(translational_vectors)):
        translational_vector = np.array(translational_vectors[i1])
        print(F'Computing the AO overlaps between R({translational_vector[0]}, {translational_vector[1]}, {translational_vector[2]}) and R(0,0,0)')
        shell_1p, l_vals =
molden_methods.molden_file_to_libint_shell(molden_filename,is_spherical,is_periodic,cell,
translational_vector)
AO_S += compute_overlaps(shell_1,shell_1p, nprocs)
```

What do we need to specify in the input?

```

params['nprocs']
params['mpi_executable']
params['istep']
params['fstep']
params['lowest_orbital']
params['highest_orbital']
params['isxTB']
params['isUKS']
params['is_periodic']
if params['is_periodic']:
    params['A_cell_vector']
    params['B_cell_vector']
    params['C_cell_vector']
    params['periodicity_type'] # example: 'XYZ'
origin = [0,0,0]
params['translational_vectors'] =
CP2K_methods.generate_translational_vectors(
    origin, [2,2,2], params['periodicity_type'])

params['is_spherical']
params['remove_molden']
params['res_dir']
params['all_pdosfiles']
params['all_logfiles']

```

```

params['cp2k_exe']
params['cp2k_ot_input_template'] # just for xTB
params['cp2k_diag_input_template']
params['trajectory_xyz_filename']

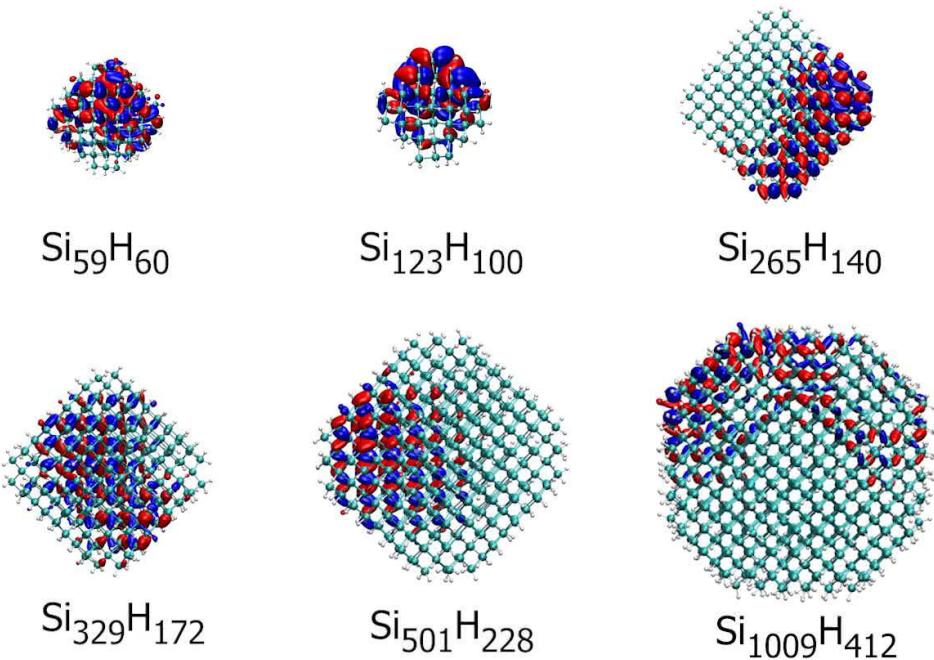
# For cube visualization
params['cube_visualization']
params['vmd_input_template']
params['states_to_plot']
params['plot_phase_corrected']
params['vmd_exe']
params['tachyon_exe']
params['x_pixels']
params['y_pixels']
params['image_format']
params['remove_cube']
params['all_images']

step2.run_cp2k_libint_step2(params)

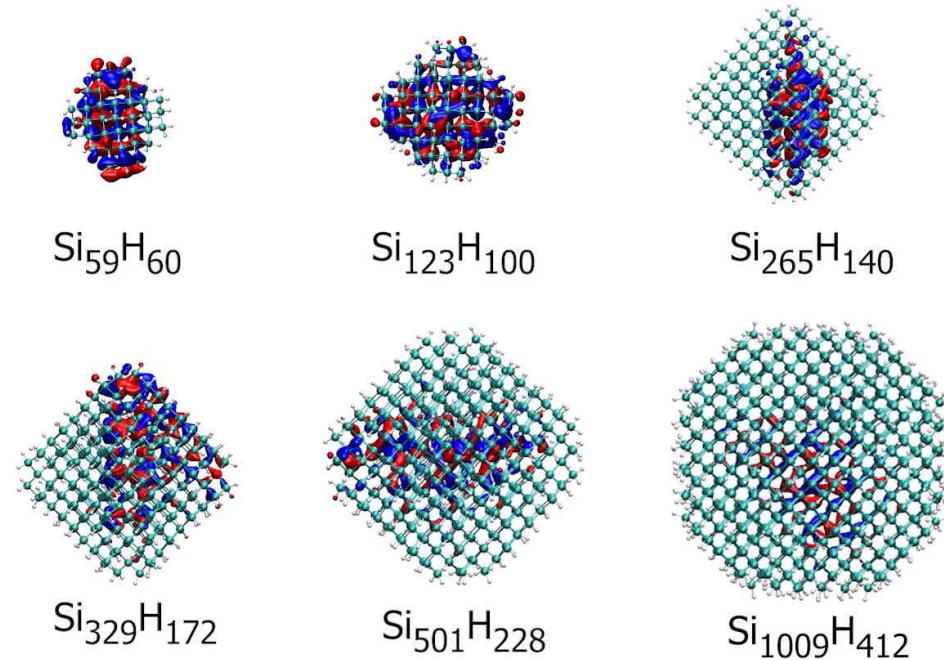
```

Molecular orbitals visualization

Highest Occupied Molecular Orbital



Lowest Unoccupied Molecular Orbital

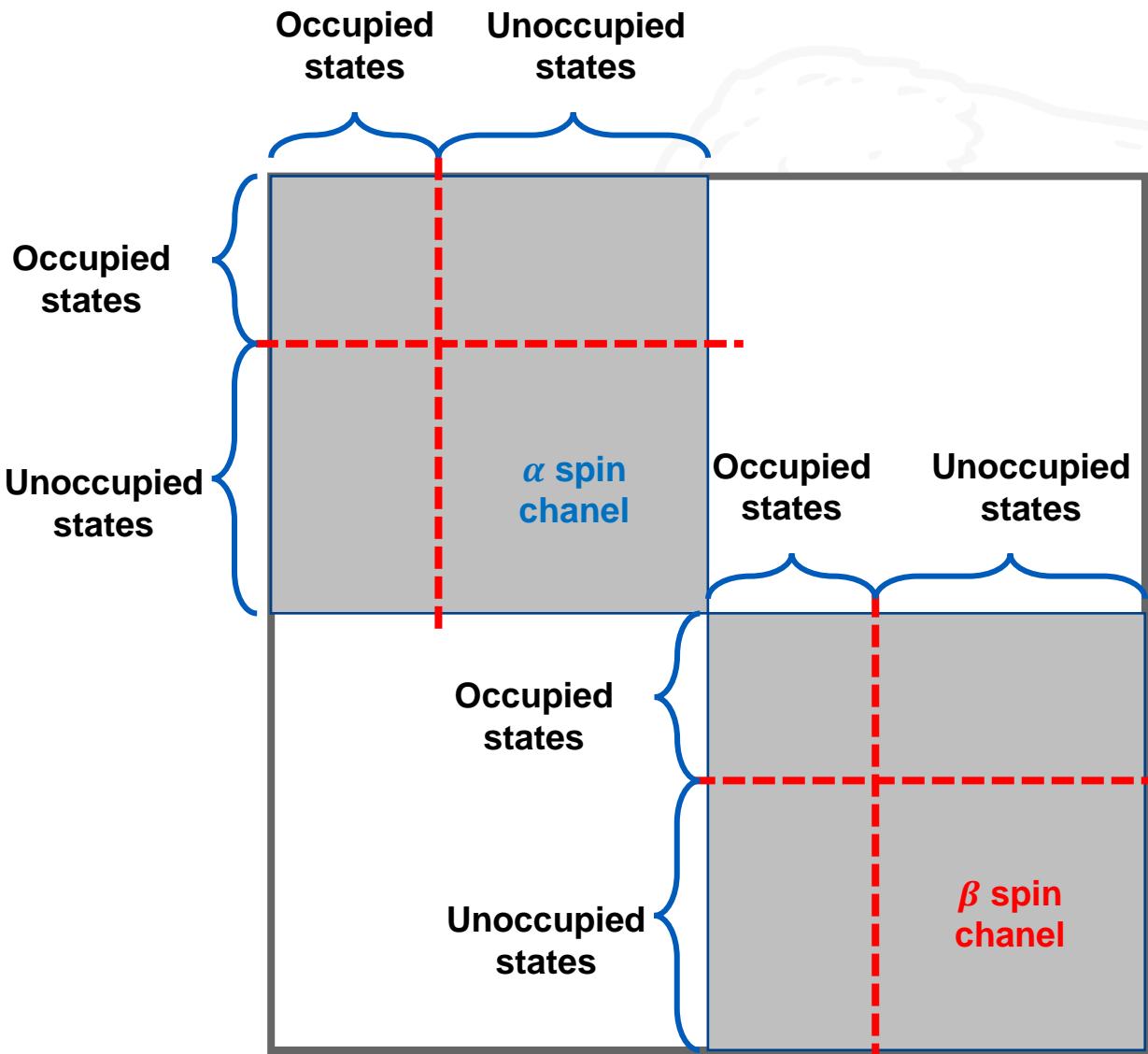


0

0

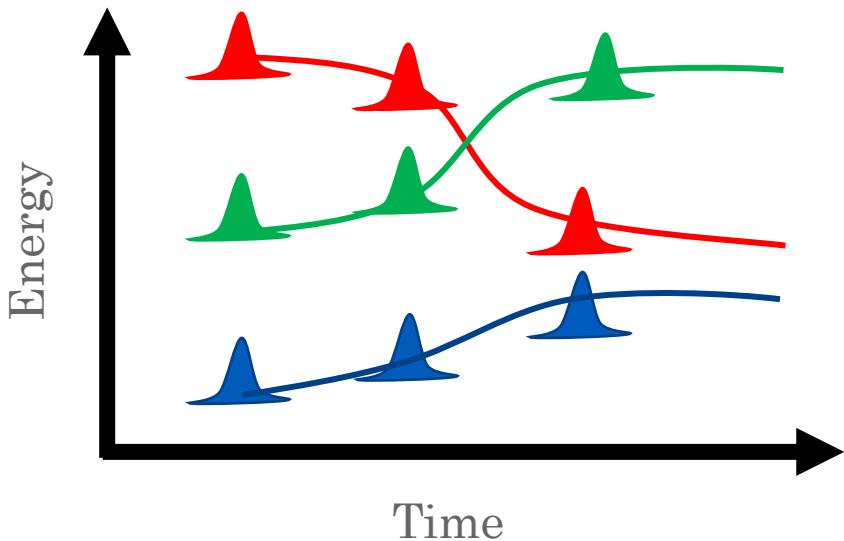
Storing data

- Libra saves MO overlaps in 2-spinor format
- With no spin-orbit coupling, the second and third block of the matrix is zero.
- To efficiently storing the data, we use `scipy.sparse` library of Python



Corrections to time-overlap matrix

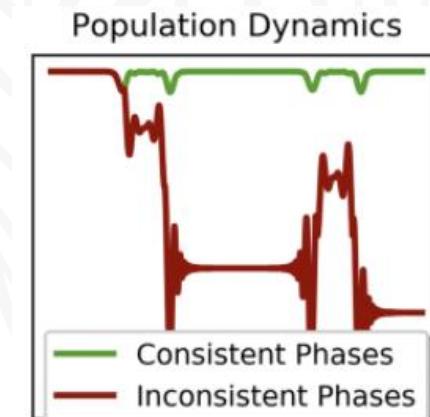
- State tracking
 - Min-cost algorithm
 - Stochastic state-tracking



- Phase correction

$$F_i(t + \Delta t) = \frac{\langle \Psi_i(t) | \Psi_i(t + \Delta t) \rangle}{\| \Psi_i(t) \|^{*} \| \Psi_i(t + \Delta t) \|}$$

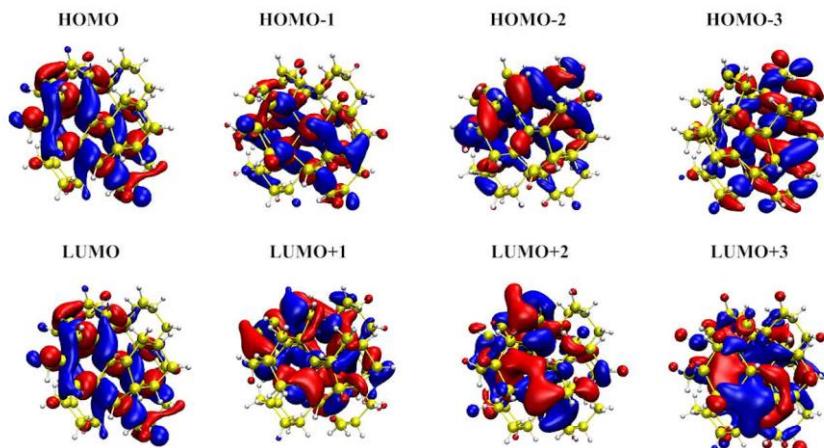
The phase-corrected wave function at time $t + \Delta t$ is:
 $\overline{\Psi}_i(t + \Delta t) = F_i^{*}(t + \Delta t) \Psi_i(t + \Delta t)$



Corrections to time-overlap matrix

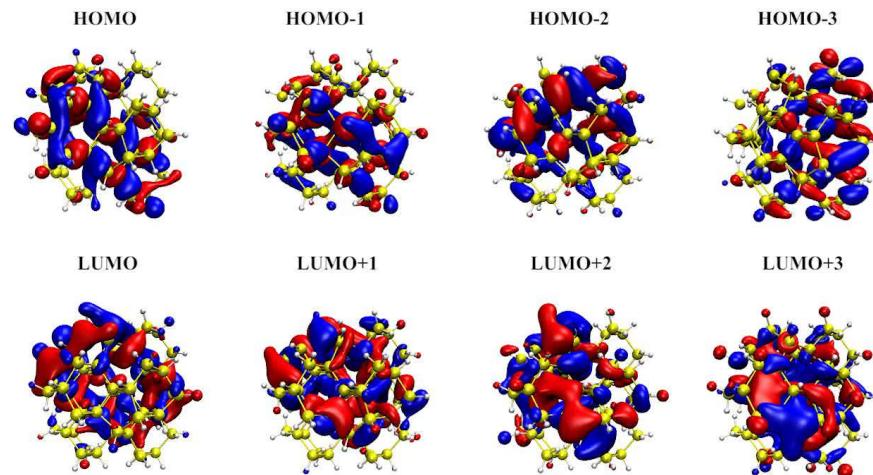
- Phase correction

Phase-uncorrected orbitals



0

Phase-corrected orbitals



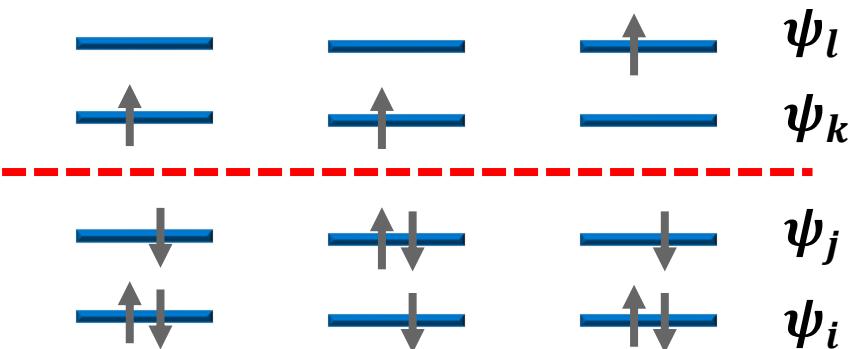
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Excited states basis

- In quantum mechanics, electronic wave function has the antisymmetric property which can be expressed in form of Slater determinant.
- For an N-electron system:

$$\Psi(x_1, \dots, x_N) = \frac{1}{\sqrt{N!}} |\psi_1(x_1)\psi_2(x_2) \dots \psi_K(x_N)\rangle = \frac{1}{\sqrt{N!}} \begin{vmatrix} \psi_1(x_1) & \cdots & \psi_K(x_1) \\ \vdots & \ddots & \vdots \\ \psi_1(x_N) & \cdots & \psi_K(x_N) \end{vmatrix}$$

- Single-particle excitations

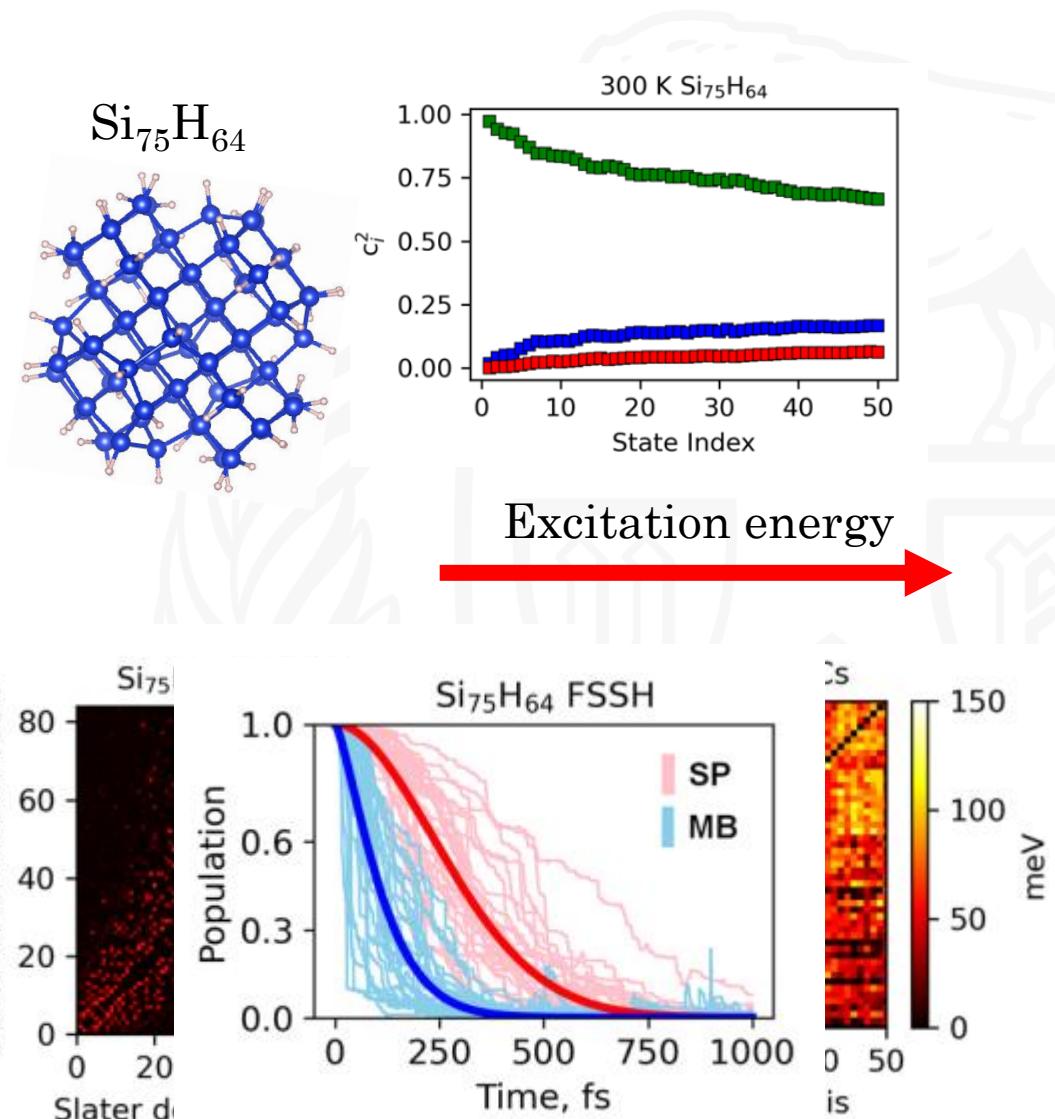
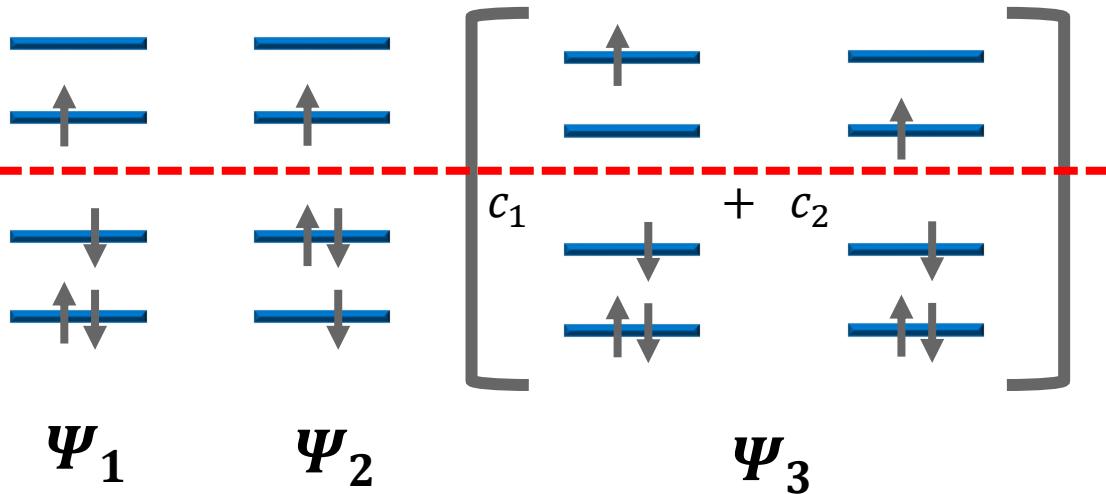


Ψ_1 Ψ_2 Ψ_3

$$\langle \Psi_1 | \Psi_2 \rangle = \det \begin{vmatrix} \langle \psi_i | \underline{\psi}_k \rangle & \langle \overline{\psi}_i | \underline{\psi}_k \rangle & \langle \underline{\psi}_i | \psi_k \rangle & \langle \overline{\psi}_j | \psi_k \rangle \\ \langle \psi_i | \underline{\psi}_i \rangle & \langle \overline{\psi}_i | \underline{\psi}_i \rangle & \langle \underline{\psi}_i | \overline{\psi}_i \rangle & \langle \overline{\psi}_j | \overline{\psi}_i \rangle \\ \langle \psi_i^2 | \overline{\psi}_j \rangle & \langle \overline{\psi}_i | \psi_j \rangle & \langle \psi_k | \psi_j \rangle & \langle \overline{\psi}_j | \psi_j \rangle \\ \langle \psi_i^3 | \overline{\psi}_j \rangle & \langle \overline{\psi}_i | \psi_l \rangle & \langle \psi_l | \psi_j \rangle & \langle \overline{\psi}_j | \overline{\psi}_j \rangle \end{vmatrix}$$

Excited states basis

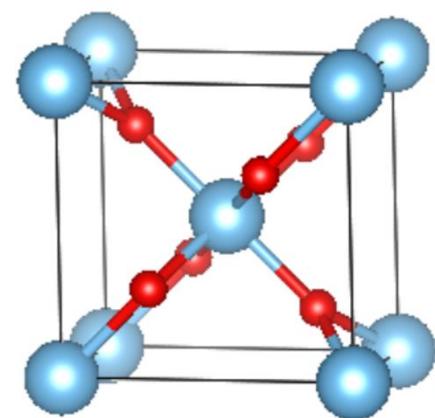
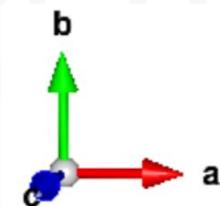
- Many-body (TD-DFT) excitation



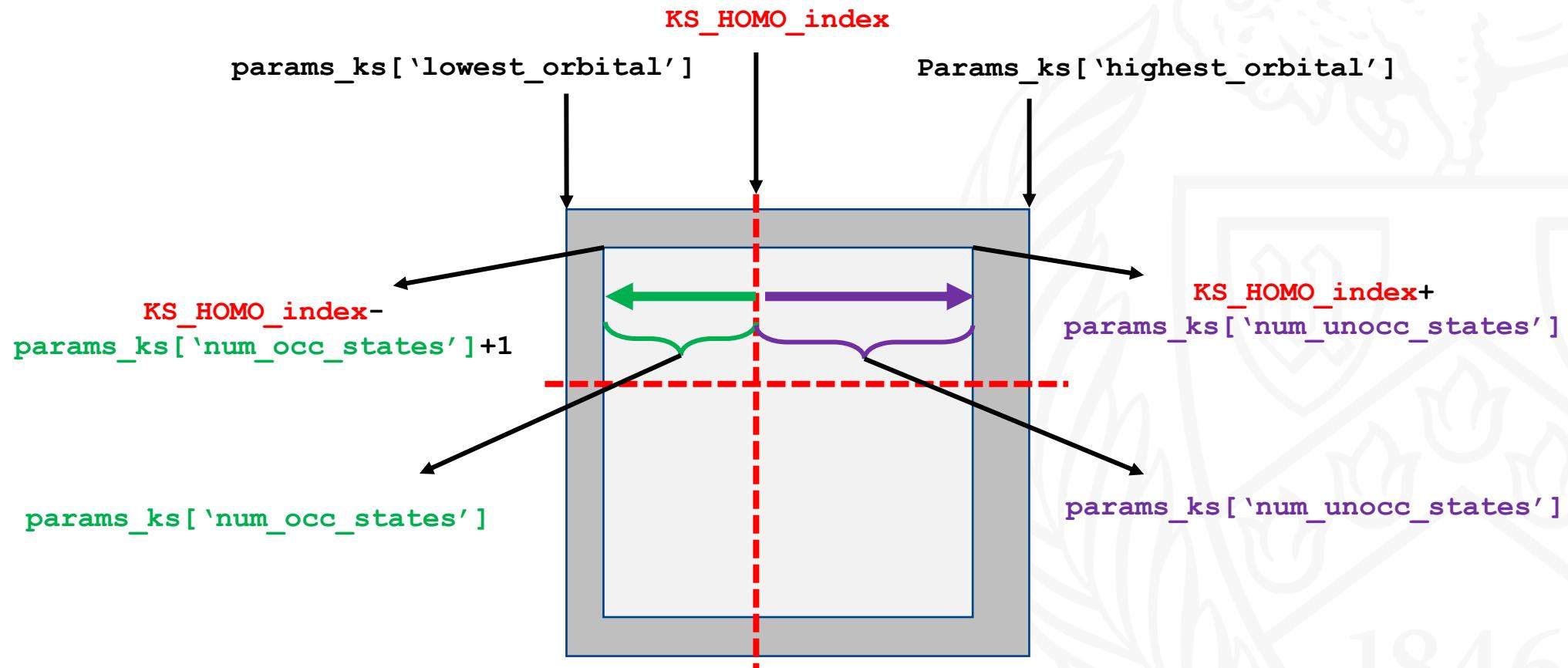
NAC in KS basis

```
params_ks = { 'lowest_orbital': 24-10, 'highest_orbital': 24+11,
  'num_occ_states': 10, 'num_unocc_states': 10, 'use_multiprocessing': True,
  'nprocs': 8, 'time_step': 1.0, 'es_software': 'cp2k', 'path_to_npz_files':
  os.getcwd()+'res', 'logfile_directory': os.getcwd()+'all_logfiles',
  'path_to_save_ks_Hvibs': os.getcwd()+'res-ks-DFT', 'start_time': 1200,
  'finish_time': 1401, 'apply_phase_correction': True,
  'apply_orthonormalization': True, 'do_state_reordering': 2,
  'state_reordering_alpha': 0, 'nac_algo': 0 }
# For KS states - Applying correction to KS overlaps and computing the
NACs in KS space
step3.run_step3_ks_nacs_libint(params_ks)
```

Rutile TiO₂
Single Unit Cell



NAC in KS basis

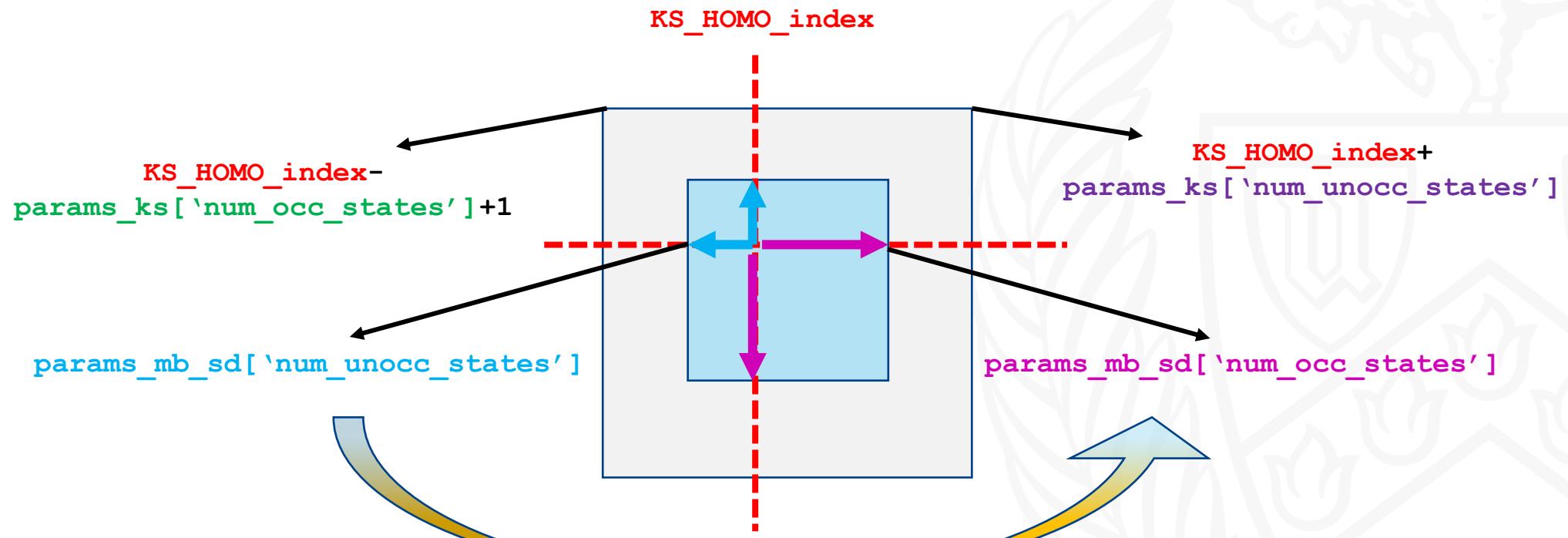


NAC in excited states basis

```
params_mb_sd = {
    'lowest_orbital': 24-params_ks['num_occ_states']+1, 'highest_orbital': 24+params_ks['num_unocc_states'],
    'num_occ_states': 10, 'num_unocc_states': 10,
    'isUKS': 0, 'number_of_states': 10, 'tolerance': 0.01, 'verbosity': 0,
    'use_multiprocessing': True, 'nprocs': 12,
    'is_many_body': True, 'time_step': 1.0, 'es_software': 'cp2k',
    #'path_to_npz_files': os.getcwd()+'//7_step2_cp2k/1_DFT/2_hpc/1_example_TiO2/res',
    'path_to_npz_files': os.getcwd()+'/res-ks-DFT',
    'logfile_directory': os.getcwd()+'//7_step2_cp2k/1_DFT/2_hpc/1_example_TiO2/all_logfiles',
    'path_to_save_sd_Hvibs': os.getcwd()+'/res-mb-sd-DFT',
    'outdir': os.getcwd()+'/res-mb-sd-DFT',
    'start_time': 1200, 'finish_time': 1401, 'sorting_type': 'identity',
    'apply_phase_correction': True, 'apply_orthonormalization': True,
    'do_state_reordering': 2, 'state_reordering_alpha':0, 'nac_algo': 0
}

step3.run_step3_sd_nacs_libint(params_mb_sd)
```

NAC in excited states basis



Summary

- All these methodologies are implemented and available in Libra software package
 - Open-source code for quantum dynamics methodologies such as trajectory surface hopping
 - The underlying code is written in C++ for faster computation and the functions can be called from Python
 - Libint is used for computation of overlaps between GTOs
 - Sparse representation of the overlap matrices using `scipy.sparse` library in Python
 - High-throughput computation for generating the overlap matrices
 - Applicable to large systems in different electronic structure calculations frameworks such as density functional tight-binding
 - It is interfaced with many quantum chemistry codes such as CP2K, Quantum ESPRESSO, and Gaussian but the Libint interface is only available for CP2K code.



Thank You!

Questions?

